

# NEUTRON SPIN STRUCTURE IN THE RESONANCE REGION AND QUARK-HADRON DUALITY

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The Thomas Jefferson National Accelerator Facility experiment E01-012 measured the  $^3\text{He}$  spin structure functions and virtual photon asymmetries in the resonance region in the range  $1.0 < Q^2 < 4.0 (\text{GeV}/c)^2$ . Our data combined with existing deep inelastic data can be used to test quark-hadron duality on  $g_1$  and  $A_1$  for  $^3\text{He}$  and the neutron. The demonstration of duality for spin structure functions will enable the use of resonance data to study nucleon spin structure in the very high  $x_{bj}$  region. Preliminary results of  $A_1^{^3\text{He}}$  will be presented as well as an overview of the experiment and theoretical developments.

## 1. Introduction

In the 70's, Bloom and Gilman<sup>1</sup> observed that the nucleon resonances average on the high  $Q^2$  scaling curve when an appropriate scaling variable is used. Since then Bloom-Gilman duality has been experimentally demonstrated for the spin independent structure function  $F_2$  of the proton and the deuteron<sup>2</sup> and for the virtual photon asymmetry  $A_1$  of the proton<sup>3</sup>.

Substantial efforts are ongoing in investigating quark-hadron duality in polarized structure functions both experimentally and theoretically. Carlson and Mukhopadhyay<sup>4</sup> showed using perturbative QCD that the structure functions in the resonance region fall with increasing  $Q^2$  at the same rate as in the deep inelastic region. The behavior of  $g_1$  in the resonance region at high  $Q^2$  is proportional to the helicity amplitude  $G_+ = g_+/Q^3$  and can be written as follow:

$$g_1 = \frac{m_N^2}{\pi m_R \Sigma_R} G_+^2 = \frac{m_N^2}{\pi m_R \Sigma_R} \frac{g_+^2}{(m_R^2 - M_N^2)^3} (1-x)^3 \quad (1)$$

where  $g_+$  is a constant and  $1/Q^2 \approx \frac{1-x}{m_R^2 - M_N^2}$  for  $X \rightarrow 1$  and  $W \approx M_R$ . In the deep inelastic region, the photon is more likely to interact with the quark having the same helicity as the nucleon. This implies that  $g_1$  and  $F_1$  behave the same way as  $x$  approaches 1 and:

$$g_1(x) \propto (1-x)^3 \text{ as } x \rightarrow 1 \quad (2)$$

Finally they predicted  $A_1$  tends to 1 as  $x \rightarrow 1$  in the scaling region and the same is true for  $A_1$  in the resonance region at high enough  $Q^2$  considering resonant and non-resonant background.

Recently, Close and Melnitchouk<sup>5</sup> studied three different conditions of SU(6) breaking applied in the resonance region under which predictions of the structure functions at large  $x$  lead to the same result as the parton model. They examined the cases where certain resonances are removed from the summation, (that is, suppression of spin- $\frac{3}{2}$ , suppression of helicity- $\frac{3}{2}$  and suppression of symmetric wave function), and found that each scenario predicts  $A_1^{n,p} \rightarrow 1$  as  $x \rightarrow 1$ .

Now that precise spin structure data in the deep inelastic region<sup>6</sup> are available, data in the resonance region is needed (especially for the neutron) in order to test duality in the polarized case. Thus the goal of the experiment E01-012 was to produce such data in the moderate  $Q^2$  region of 1.0 to 4.0(GeV/c)<sup>2</sup> where duality is expected to hold.

## 2. The E01-012 experiment

E01-012 ran successfully in January-February 2003 at Jefferson Lab in Hall A (see fig. 1). It was an inclusive experiment of longitudinally polarized electrons scattering on a longitudinally or transversely polarized <sup>3</sup>He target<sup>7</sup>. Asymmetries and cross section differences are formed in order to extract the spin structure function  $g_1$  and the virtual photon asymmetry  $A_1$  in the resonance region up to  $Q^2 = 4(\text{GeV}/c)^2$ :

$$g_1 = \frac{MQ^2\nu}{4\alpha_e^2} \frac{E}{E'} \frac{1}{E+E'} \left[ \Delta\sigma_{\parallel} + \tan\left(\frac{\theta}{2}\right)\Delta\sigma_{\perp} \right] \quad (3)$$

$$A_1 = \frac{A_{\parallel}}{D(1+\eta\xi)} - \frac{\eta A_{\perp}}{d(1+\eta\xi)} \quad (4)$$

where  $A_{\parallel}$  ( $A_{\perp}$ ) is the parallel (perpendicular) asymmetry corrected for data

acquisition deadtime, beam charge asymmetry, target and beam polarizations and nitrogen dilution.  $\Delta\sigma_{\parallel(\perp)} = 2A_{\parallel(\perp)}\sigma_0$  with  $\sigma_0$  the unpolarized cross section, and  $\eta$ ,  $\xi$ ,  $D$  and  $d$  are kinematic factors (see for example <sup>6</sup>). To determine  $D$  and  $d$ , world data for  $R(x, Q^2)$  will be used. However our data allows a direct extraction of  $g_1$  and  $g_2$  without the need of external input.

The beam polarization measured using a Moller polarimeter was on average  $77(1 \pm 0.03)\%$ . The polarization of the target was determined by two independent polarimetries<sup>7</sup>: NMR and EPR. From the preliminary analysis, the target polarization was  $37(1 \pm 0.04)\%$  on average. Two almost identical spectrometers were used in a symmetric configuration in order to double our statistics and check our systematics. In order to select good scattered electrons, a gas cerenkov counter along with a two-layer electromagnetic calorimeter was used in the analysis. These allowed us to reduce the pion contamination by a factor better than  $10^4$  while keeping the electron efficiency above 99%.

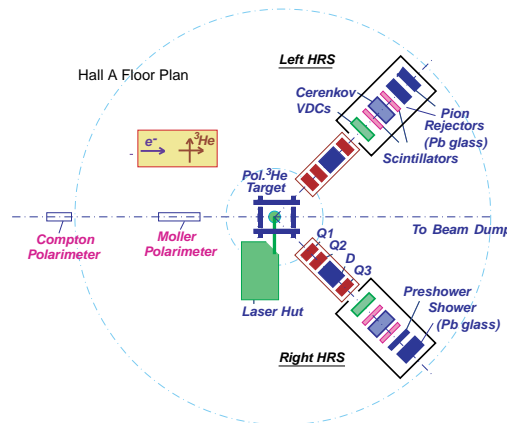


Figure 1. Hall A floor plan. The electron beam is coming from the left.

### 3. Preliminary results

For our first pass analysis,  $A_1^{3\text{He}}$  was extracted from the asymmetries.  $R(x, Q^2)$  was assumed constant with a value of 0.18. No radiative corrections are applied and the error bars are statistical only. Figure 2 shows  $A_1^{3\text{He}}$  at four different  $Q^2$  values in function of the Nachtmann scaling vari-

able  $\xi = 2x/(1 + \sqrt{1 + \frac{4M^2 x^2}{Q^2}})$ . The position of the  $\Delta(1232)$  resonance is indicated in each plot.

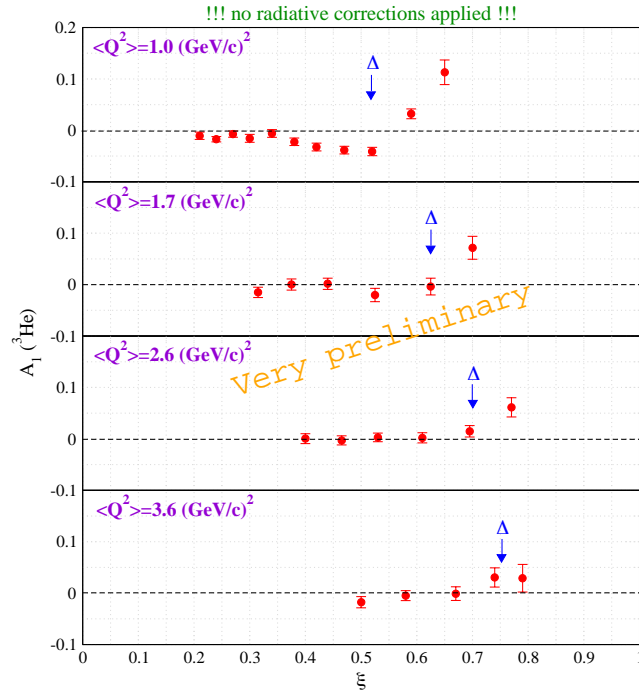


Figure 2. Preliminary result of  $A_1^{^3\text{He}}$ . See text for discussions.

The most important feature is the negative contribution of the  $\Delta(1232)$  resonance at low  $Q^2$ . It is reported <sup>4,5</sup> that quark-hadron duality is not expected to apply in the  $\Delta$  region at this low  $Q^2$ . However at higher  $Q^2$ ,  $A_1^{^3\text{He}}$  in the  $\Delta(1232)$  increases due to the fall off of the resonance and the rising background. At large  $\xi$ , our data tend to follow the same pattern as the DIS world data<sup>8</sup> and indicate the validity of duality.

#### 4. Summary

E01-012 resonance data cover the region of  $0.2 < x < 0.90$  (see fig. 3). At  $x < 0.60$  where DIS data are available, we will provide a precision test of quark-hadron duality predictions for neutron spin structure functions. Moreover, if duality is confirmed, E01-012 will provide the first precise measure-

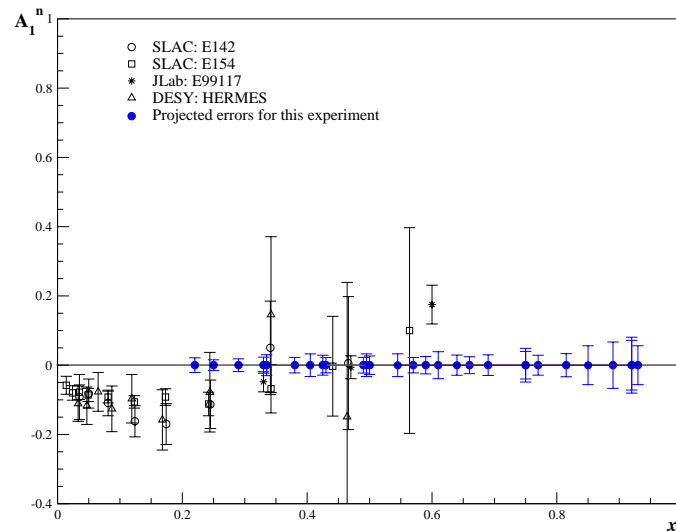


Figure 3.  $A_1^n$  results measured in the deep inelastic region using polarized  $^3\text{He}$  targets and E01-012 projected statistical errors are shown.

ment of  $g_1^n$  and  $A_1^n$  in the range  $0.60 < x < 0.90$ .

Our data will also be used to extract moments of the structure functions, for example: the extended GDH sum,  $d_2$  matrix element and the Burkhardt-Cottingham sum rule.

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